Wavelet Transforms and Symmetric Tube Domains

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Abstract. We extend wavelet analysis to the symmetric tube domains and their Shilov boundaries. Our approach is based on the theory of Jordan algebra.

One-dimensional wavelet analysis can be explained in terms of square-integrable representation of the affine group (cf. [4], [6]). It is an intermediate between the function theory on the upper half-plane of one complex variable and the harmonic analysis on the real line (cf. [7], [9]). In this paper we extend wavelet analysis to the symmetric tube domains and their Shilov boundaries, the higher dimensional analogues of the upper half-plane and the real line. We assume that V is a simple Euclidean Jordan algebra, Ω is the associated symmetric cone and T_{Ω} is the symmetric tube domain over Ω . In §1, we recall some notations and facts about Jordan algebras and symmetric cones, especially the Iwasawa subgroup P of the holomorphic automorphism group of T_{Ω} . P has a natural unitary representation π on $L^2(V)$. In §2, we decompose $L^2(V)$ into the direct sum of the irreducible invariant closed subspaces under π . In §3, we give an explicit characterization of the admissibility condition in terms of Fourier transform and Jordan algebra. We also give a family of admissible wavelets, which is a complete orthonormal system in a sense. Finally in §4, we use wavelet transforms to decompose the weighted L^2 -space on the tube domain T_{Ω} into a direct sum of subspaces such that the first component is exactly the weighted Bergman space.

A good reference on Jordan algebras, symmetric cones and tube domains is the book [3] by J. Faraut and A. Korányi. Various authors developed the theory of continuous wavelet in view of square-integrable group representations, for example, in [5], [8] and in particular [1].

1. Iwasawa subgroup

Throughout this paper we keep the following assumptions and notations, which are the same as in [3].

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V is an *n*-dimensional simple Euclidean Jordan algebra with identity e. xy denotes the Jordan product of x and y. tr(x) and det(x) are defined as in [3]. We also write $\Delta(x)$ instead of det(x). The inner product on V is given by (x|y) = tr(xy). L(x) is the linear map of V defined by L(x)y = xy. An element $c \in V$ is idempotent if $c^2 = c$. The only eigenvalues of L(c) are $1, \frac{1}{2}$, and 0. The corresponding eigenspaces are denoted by V(c, 1), $V(c, \frac{1}{2})$ and V(c, 0). We fix a Jordan frame $\{c_1, \dots, c_r\}$, where r is the rank of V. Then we have the Peirce decomposition

$$V = \bigoplus_{j \le k} V_{jk}$$

where

$$V_{ii} = V(c_i, 1) = \mathbf{R}c_i,$$

 $V_{ij} = V(c_i, \frac{1}{2}) \cap V(c_j, \frac{1}{2}).$

 $d = \dim V_{ij}$, which does not depend on i and j, is called the degree of V. Let

$$P(x) = 2L(x)^2 - L(x^2)$$

be the quadratic representation, and write

$$x\Box y = L(xy) - [L(x), L(y)]$$

For given j and for $z^{(j)} \in \bigoplus_{k=j+1}^{r} V_{jk}$ the Frobenius transform $\tau(z^{(j)})$ is defined by

$$\tau(z^{(j)}) = \exp(2z^{(j)} \Box c_j).$$

Let Ω be the symmetric cone which consists of elements x in V such that L(x) is positive definite. $G(\Omega)$ denotes the automorphism group of Ω and G is the identity component of $G(\Omega)$. G has Iwasawa decomposition G = NAK, where

$$K = \{g \in G : ge = e\},\$$

$$A = \{P(a) : a = \sum_{j=1}^{r} a_j c_j, a_j > 0\},\$$

$$N = \{\tau(z^{(1)}) \cdots \tau(z^{(r-1)}) : z^{(j)} \in \bigoplus_{k=j+1}^{r} V_{jk}\}\$$

are compact, diagonal and strict triangular respectively. A normalizes N and

(1.1) $P(a)\tau(z^{(j)}) = \tau(\tilde{z}^{(j)})P(a)$

where

$$z^{(j)} = \sum_{j < k} z_{jk}, \quad z_{jk} \in V_{jk},$$
$$\tilde{z}^{(j)} = \sum_{j < k} \tilde{z}_{jk}, \quad \tilde{z}_{jk} \in V_{jk},$$
$$\tilde{z}_{jk} = \frac{a_k}{a_j} z_{jk}.$$

T=NA is a semi-direct product. We will use another parametrization of the triangular subgroup $T.\ {\rm Set}$

$$V_{+} = \{ u = \sum_{j=1}^{r} u_{j}c_{j} + \sum_{j < k} u_{jk} : u_{j} > 0, u_{jk} \in V_{jk} \}.$$

For $u \in V_+$, we define

$$t(u) = P(b_1)\tau(u^{(1)})P(b_2)\cdots\tau(u^{(r-1)})P(b_r),$$

where

$$b_j = c_1 + \dots + c_{j-1} + u_j c_j + c_{j+1} + \dots + c_r,$$

 $u^{(j)} = \sum_{k=j+1}^r u_{jk}.$

Then

$$T = \{t(u) : u \in V_+\}.$$

Using (1.1), it is easy to determine the left and right Haar measures of T. The left Haar measure of T is given by

$$d\mu_l(t(u)) = 2^r \prod_{j=1}^r u_j^{-d(j-1)-1} du,$$

and the right Haar measure of T is given by

$$d\mu_r(t(u)) = 2^r \prod_{j=1}^r u_j^{-d(r-j)-1} du.$$

T acts simply and transitively on $\Omega.$ If

$$x = \sum_{j=1}^{r} x_j c_j + \sum_{j < k} x_{jk}$$

is the Peirce decomposition of x = t(u)e, then

$$x_j = u_j^2 + \frac{1}{2} \sum_{k=1}^{j-1} ||u_{kj}||^2,$$
$$x_{jk} = u_j u_{jk} + 2 \sum_{l=1}^{j-1} u_{lj} u_{lk}.$$

We identify Ω with T by identification of x = t(u)e and t(u). Then we have

$$dx = 2^r \prod_{j=1}^r u_j^{d(r-j)+1} du,$$
$$\Delta(x) = \prod_{j=1}^r u_j^2.$$

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Therefore

$$\Delta(x)^{-\frac{n}{r}}dx = d\mu_l(t(u))$$

which gives the G-invariant measure on Ω .

Let V^C denote the complexification of V. $T_{\Omega} = V + i\Omega$ is the tube domain over Ω in V^C . $G(T_{\Omega})$ denotes the holomorphic automorphism group of T_{Ω} and $G(T_{\Omega})^0$ is the identity component of $G(T_{\Omega})$. The Iwasawa decomposition of $G(T_{\Omega})^0$ is given by $G(T_{\Omega})^0 = \underline{N}A\underline{K}$, where

$$\underline{K} = \{g \in G(T_{\Omega})^0 : g(ie) = ie\} \supset K,$$

$$\underline{N} = N^+ N,$$

$$N^+ = \{\tau_u : z \mapsto z + u, \quad u \in V\} \cong V.$$

Therefore,

$$G(T_{\Omega})^0 = N^+ T \underline{K}.$$

We call it the partial Iwasawa decomposition as in Terras' book [11]. T normalizes N^+ as

$$t(v)\tau_u = \tau_{t(v)u}t(v), \quad u \in V, \, v \in V_+$$

 $P = \underline{N}A = N^+T$ is called the Iwasawa subgroup. P is a nonunimodular group. Using the parametrization (u, v) for $\tau_u t(v) \in P$, the left Haar measure of P is given by

$$d\mu_l(u,v) = 2^r \prod_{j=1}^r v_j^{-d(r+j-2)-3} du \, dv = \prod_{j=1}^r v_j^{-d(r-1)-2} du \, d\mu_l(t(v)),$$

and the right Haar measure of P is given by

$$d\mu_r(u,v) = 2^r \prod_{j=1}^r v_j^{-d(r-j)-1} du \, dv = du \, d\mu_r(t(v)).$$

P acts on T_{Ω} simply and transitively. We identify T_{Ω} with P by identification of $\tau_u t(v)(ie)$ and $\tau_u t(v)$. If $x + iy = \tau_u t(v)(ie) = u + it(v)e$, then

$$\Delta(y)^{-\frac{2n}{r}}dx\,dy = d\mu_l(u,v),$$

which is the $G(T_{\Omega})^0$ -invariant measure on T_{Ω} . Note that

$$\operatorname{Det}(g) = \Delta(ge)^{\frac{n}{r}}, \quad g \in G.$$

P has a natural unitary representation on $L^2(V)$ defined by

$$\pi_{(u,v)}: f(x) \mapsto \Delta(t(v)e)^{-\frac{n}{2r}} f(t(v)^{-1}x - t(v)^{-1}u).$$

We shall decompose $L^2(V)$ into the direct sum of irreducible invariant closed subspaces under π .

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2. The decomposition of $L^2(V)$

In order to decompose $L^2(V)$, we need to identify the non-degenerate *T*-orbits of *V* under the contragredient action of *T*, which is given by $x \mapsto t(v)'^{-1}x$ where t(v)' denotes the transpose of t(v). First we prove

Lemma 1. (1) Suppose $z_{ij} \in V_{ij}$, $w_{kl} \in V_{kl}$, i < j, k < l, $i \neq l$, $k \neq j$, then

$$[z_{ij}\Box c_i, w_{kl}\Box c_k] = 0.$$

(2) Suppose $z_{ij} \in V_{ij}$, then

$$(z_{ij} \Box c_i)' = z_{ij} \Box c_j.$$

Proof. (a) To prove (1), we use the facts

$$V_{ij} \cdot V_{jk} \subset V_{ik}, \quad \text{if} \quad i \neq k,$$

$$V_{ij} \cdot V_{kl} = \{0\}, \quad \text{if} \quad \{i, j\} \cap \{k, l\} = \emptyset,$$

$$xy = \frac{1}{2}(x|y)(c_i + c_j), \quad \text{if} \quad x, y \in V_{ij}$$

(cf [3], Theorem IV.2.1 (iii) and Proposition IV.1.4 (i)). We also use the matrix of $z \Box c$ with respect to the Peirce decomposition, when c is idempotent in V and $z \in V(c, \frac{1}{2})$ (see [3], proof of Lemma VI.3.1). Let

$$x = \sum_{j=1}^{r} x_j c_j + \sum_{j < k} x_{jk}, \quad x_{jk} \in V_{jk}.$$

We compute separately in four cases.

1) If $k = i, \ l = j, \ i < j$, then

$$(z_{ij} \Box c_i)(w_{ij} \Box c_i)x = \frac{x_i}{4}(z_{ij}|w_{ij})c_j = (w_{ij} \Box c_i)(z_{ij} \Box c_i)x.$$

2) If $k = i, l \neq j, i < j, l$, then

$$(z_{ij}\Box c_i)(w_{il}\Box c_i)x = \frac{x_i}{2}z_{ij}w_{il} = (w_{il}\Box c_i)(z_{ij}\Box c_i)x.$$

3) If
$$k \neq i$$
, $l = j$, $i, k < j$, then

$$(z_{ij}\Box c_i)(w_{kj}\Box c_k)x = \frac{1}{2}(z_{ij}|x_{ik}w_{kj})c_j = \frac{1}{2}(w_{kj}|x_{ik}z_{ij})c_j = (w_{kj}\Box c_k)(z_{ij}\Box c_i)x,$$

where the second equality is due to the associativity of the inner product.

4) If $k \neq i, j, \ l \neq i, j, \ i < j, \ k < l$, we may assume i < k, then

$$(z_{ij}\Box c_i)(w_{kl}\Box c_k)x = z_{ij}(x_{ik}w_{kl}) = w_{kl}(x_{ik}z_{ij}) = (w_{kl}\Box c_k)(z_{ij}\Box c_i)x,$$

where the second equality follows from the Lemma V.3.2 in [3].

(b) Take $x = z_{ij}, y = c_i + c_j$ in the identity

$$[L(x), L(y^{2})] + 2[L(y), L(xy)] = 0$$

(cf [3]. Proposition II.1.1), we obtain

$$[L(c_i), L(z_{ij})] = [L(z_{ij}), L(c_j)].$$

It follows that

$$(z_{ij}\Box c_i)' = c_i\Box z_{ij} = z_{ij}\Box c_j$$

Let $z_{jk} \in V_{jk} (j < k)$ and put

$$z^{(j)} = \sum_{k=j+1}^{r} z_{jk}, \quad z_{(k)} = \sum_{j=1}^{k-1} z_{jk}.$$

Put

$$\tau'(z_{(k)}) = \exp(2z_{(k)}\Box c_k).$$

If $z_{ij} \in V_{ij}$, $w_{kl} \in V_{kl}$, i < j, k < l, $i \neq l$, $k \neq j$, Lemma 1 implies that

$$\tau(z_{ij})\tau(w_{kl}) = \tau(w_{kl})\tau(z_{ij})$$

and

$$\tau(z_{ij})' = \tau'(z_{ij}).$$

Thus $\tau'(z_{ij})$ is a dual Frobenius transform. Also, by Lemma 1,

$$\tau(z^{(j)}) = \tau(z_{j,j+1}) \cdots \tau(z_{j,r}), \tau'(z_{(k)}) = \tau'(z_{1,k}) \cdots \tau'(z_{k-1,k}).$$

Therefore, for

$$u = \sum_{j=1}^{r} u_j c_j + \sum_{j < k} u_{jk}, \quad u_j > 0, \, u_{jk} \in V_{jk},$$

we have, by also using (1.1),

$$t(u) = P(b_1)\tau(u^{(1)})P(b_2)\cdots\tau(u^{(r-1)})P(b_r)$$

= $P(b_1)\tau(u_{12})P(b_2)\tau(u_{13})\tau(u_{23})\cdots P(b_{r-1})\tau(u_{1r})\cdots\tau(u_{r-1,r})P(b_r).$
 $t(u)' = P(b_r)\tau'(u_{r-1,r})\tau'(u_{r-2,r})\cdots\tau'(u_{1r})P(b_{r-1})\cdots P(b_2)\tau'(u_{12})P(b_1)$
= $P(b_r)\tau'(u_{(r)})P(b_{r-1})\cdots\tau'(u_{(2)})P(b_1)$

where

$$u_{(k)} = \sum_{j=1}^{k-1} u_{jk}.$$

For $j = 1, \dots, r$, let $V^{(j)}$ be the subalgebra $V(c_1 + \dots + c_j, 1)$ of V and $W^{(j)}$ be the subalgebra $V(c_{r-j+1} + \dots + c_r, 1)$ of V. P_j and P_j^* denote the orthogonal projections onto $V^{(j)}$ and $W^{(j)}$ respectively. $\det_{(j)}$ and $\det_{(j)}^*$ are the determinants relative to $V^{(j)}$ and $W^{(j)}$ respectively. We define

$$\Delta_j(x) = \det_{(j)}(P_j x),$$

$$\Delta_j^*(x) = \det_{(j)}^*(P_j^* x).$$

Furthermore, for $\mathbf{s} = (s_1, \cdots, s_r)$. We let

$$\Delta_{\mathbf{s}}(x) = \Delta_{1}(x)^{s_{1}-s_{2}} \cdots \Delta_{r-1}(x)^{s_{r-1}-s_{r}} \Delta_{r}(x)^{s_{r}}, \Delta_{\mathbf{s}}^{*}(x) = \Delta_{1}^{*}(x)^{s_{1}-s_{2}} \cdots \Delta_{r-1}^{*}(x)^{s_{r-1}-s_{r}} \Delta_{r}^{*}(x)^{s_{r}}.$$

For $x \in V, t(u) \in T$, we have

(2.1)
$$\Delta_{\mathbf{s}}^*(t(u)'x) = u_1^{2s_r} \cdots u_r^{2s_1} \Delta_{\mathbf{s}}^*(x) = \Delta_{\mathbf{s}}^*(t(u)'e) \Delta_{\mathbf{s}}^*(x).$$

In particular, $\Delta_{\mathbf{s}}^*$ is invariant under the Frobenius transform $\tau'(z_{(k)})$ (cf [3], Proposition VII.1.5).

Set

$$E = \{ \varepsilon = \sum_{j=1}^{r} \varepsilon_j c_j : \varepsilon_j = 1 \text{ or } i \},\$$
$$\Omega_{\varepsilon} = \{ x \in V : x = t(u)' P(\varepsilon) e, u \in V_+ \}$$

Lemma 2. (1) The Ω_{ε} 's are disjoint and simply transitive orbits under the contragredient action of T. (2) $\bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$ is a set with a complementary of measure zero.

Proof. (a) Suppose that

$$t(u)'P(\varepsilon)e = t(v)'P(\delta)e, \quad u, v \in V_+, \varepsilon, \delta \in E$$

Write

$$g = P(\delta)t(v)'^{-1}t(u)'P(\varepsilon).$$

Since t(u), t(v) are triangular and $P(\varepsilon), P(\delta)$ are diagonal, g is triangular. On the other hand, since ge = e, from the Proposition VIII.2.4 in [3] g is an automorphism of V^C and $g' = g^{-1}$. Therefore g is diagonal. Because t(u), t(v)have positive diagonal elements and $P(\varepsilon), P(\delta)$ have diagonal elements 1, -1 or i, it is concluded that $u = v, \varepsilon = \delta$.

(b) Set

$$B = \{x \in V : \Delta_k^*(x) \neq 0, k = 1, \cdots, r\}$$

Obviously, $V \setminus B$ is a zero measure set. We will prove that $B = \bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$. It is easy to see that $B \supset \bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$. Assume that

$$x = \sum_{j=1}^{r} x_j c_j + \sum_{j < k} x_{jk} \in B.$$

By [3], Theorem VI.3.5 we can write

$$x = \tau'(z_{(r)}) \cdots \tau'(z_{(2)}) \sum_{j=1}^{r} a_j c_j$$

where

$$z_{(k)} = \sum_{j=1}^{k-1} z_{jk} \in \bigoplus_{j=1}^{k-1} V_{jk},$$

$$a_j = \frac{\Delta_{r-j+1}^*(x)}{\Delta_{r-j}^*(x)} \neq 0, \quad j = 1, \cdots, r-1,$$

$$a_r = \Delta_1^*(x) \neq 0.$$

Set

$$\varepsilon_j = \begin{cases} 1, & \text{if } a_j > 0, \\ i, & \text{if } a_j < 0, \end{cases}$$
$$u_j = \sqrt{|a_j|},$$
$$u_{jk} = u_k z_{jk}.$$

Then, by (1.1),

 $x = t(u)' P(\varepsilon)e.$

Remark. Clearly, $\Omega_e = \Omega$, $\Omega_{ie} = -\Omega$. Ω_{ε} is a connected open set in V because Ω_{ε} is homeomorphic to V_+ . But Ω_{ε} may not be convex neither K-invariant in general.

A simple example of Lemma 2 can be given as follows. Let V be the space $\operatorname{Sym}(m, \mathbf{R})$ of all $m \times m$ symmetric matrices and $c_j = \operatorname{diag}(0, \dots, 0, 1, 0, \dots, 0)$. An element t in T has the following form: tx = uxu', where u is a lower triangular matrix with positive diagonal elements. Let Σ denote the set of all diagonal matrices with diagonal elements ± 1 . $\Omega_{\sigma}(\sigma \in \Sigma)$ consists of all matrices of form $u'\sigma u$. Then Ω_{σ} 's are disjoint and simply transitive orbits under the adjoint action of T and $\bigcup_{\sigma \in \Sigma} \Omega_{\sigma}$ is a total measure set. Now we are ready to decompose $L^2(V)$. Set

$$H_{\varepsilon} = \{ f \in L^2(V) : \operatorname{supp} \tilde{f} \subseteq \operatorname{Cl}(\Omega_{\varepsilon}) \}.$$

Proposition 1. Each of H_{ε} is an irreducible invariant closed subspace of $L^{2}(V)$ under π and

(2.2)
$$L^2(V) = \bigoplus_{\varepsilon \in E} H_{\varepsilon}.$$

Proof. (2.2) follows from Lemma 2. Because

$$(\pi_{(u,v)}f)(y) = \Delta(t(v)e)^{\frac{n}{2r}}e^{-i(u|y)}\hat{f}(t(v)'y),$$

it is easy to see that H_{ε} is invariant under π . We need to prove that H_{ε} is irreducible. Let W be a non-zero invariant closed subspace of H_{ε} under π and W^+ the orthogonal complement of W in H_{ε} . Taking a function $g \in W$, not identically zero, if $f \in W^+$, then

$$\langle f, \pi_{(u,v)}g \rangle_{L^2(V)} = \int_V f(x)\overline{\pi_{(u,v)}g(x)} \, dx = 0, \quad u \in V, \, v \in V_+.$$

Write

$$\tilde{g}(x) = \overline{g(-x)},$$

$$g_{t(v)}(x) = \Delta(t(v)e)^{-\frac{n}{2r}}g(t(v)^{-1}x).$$

We have

(2.3)
$$\langle f, \pi_{(u,v)}g \rangle_{L^2(V)} = f * \tilde{g}_{t(v)}(u).$$

Therefore,

(2.4)
$$(f * \tilde{g}_{t(v)})^{\hat{}}(y) = \Delta(t(v)e)^{\frac{n}{2r}} \hat{f}(y)\overline{\hat{g}(t(v)'y)} = 0, \quad a.e. \ y \in V.$$

Set

$$S_1 = \operatorname{supp} \hat{f} \cap \Omega_{\varepsilon},$$

$$S_2 = \operatorname{supp} \hat{g} \cap \Omega_{\varepsilon}.$$

 S_1^d and S_2^d consist of points of density of S_1 and S_2 respectively. S_2^d is a positive measure set since g is not identically zero. If S_1^d has positive measure, by Lemma 2, there exists $t(v_0) \in T$ such that $S = S_1^d \cap t(v_0)'^{-1}S_2^d$ has positive measure. But

 $(f * \tilde{g}_{t(v)})(y) \neq 0, \quad y \in S,$

which contradicts (2.4). Therefore f is identically zero. This proves that H_{ε} is irreducible.

Remark. For F in $H^2(T_{\Omega})$, the Hardy space on T_{Ω} , the following limit exists,

$$\lim_{y \to 0, y \in \Omega} F(\cdot + iy) = f, \quad \text{in } L^2(V).$$

Then

$$H^2(V) = \{ f \in L^2(V) : \text{ there exists } F \in H^2(T_\Omega) \text{ such that } f = \lim F \}$$

is called the Hardy space on V. It is easy to see that $H_e = H^2(V)$ and $H_{ie} = \overline{H^2(V)}$.

3. The admissibility condition

The restriction of π on H_{ε} is square-integrable, *i.e.*, there exists a function $\phi(\neq 0)$ in H_{ε} such that

(3.1)
$$C_{\phi} = \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{P} |\langle \phi, \pi_{(u,v)}\phi \rangle_{L^{2}(V)}|^{2} d\mu_{l}(u,v) < \infty.$$

(3.1) is called the admissibility condition and ϕ is called an admissible wavelet. We want to give a characterization of the admissibility condition in terms of Fourier transform and Jordan algebra, which does not involve any group representation.

Lemma 3. Suppose $x = t(u)'P(\varepsilon)e$ in Ω_{ε} . If

$$x = \sum_{j=1}^{r} x_j c_j + \sum_{j < k} x_{jk}$$

is the Peirce decomposition of x, then

$$x_{j} = \varepsilon_{j}^{2} u_{j}^{2} + \frac{1}{2} \sum_{k=j+1}^{r} \varepsilon_{k}^{2} ||u_{jk}||^{2},$$
$$x_{jk} = \varepsilon_{k}^{2} u_{k} u_{jk} + 2 \sum_{l=k+1}^{r} \varepsilon_{l}^{2} u_{jl} u_{kl}.$$

Lemma 3 can be proved in a similar way as in [3], Proposition VI.3.8.

For the transformation $x = t(u)' P(\varepsilon)e$, by Lemma 3, it is easy to compute that

$$dx = 2^{r} \prod_{j=1}^{r} u_{j}^{d(j-1)+1} du$$
$$= \prod_{j=1}^{r} u_{j}^{2d(j-1)+2} d\mu_{l}(t(u)).$$

Let

$$\underline{\mathbf{s}} = (1 + d(r-1), 1 + d(r-2), \cdots, 1)$$

By (2.1),

$$\Delta_{\underline{\mathbf{s}}}^*(x) = \Delta_{\underline{\mathbf{s}}}^*(t(u)'e)\Delta_{\underline{\mathbf{s}}}^*(P(\varepsilon)e)$$

Therefore,

$$|\Delta_{\underline{\mathbf{s}}}^*(x)| = \Delta_{\underline{\mathbf{s}}}^*(t(u)'e).$$

and we have

(3.2)
$$|\Delta_{\underline{\mathbf{s}}}^*(x)|^{-1} dx = d\mu_l(t(u)).$$

We denoted by AW_{ε} the set of all admissible wavelets in H_{ε} .

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$$C_{\phi} = \int_{\Omega_{\varepsilon}} |\hat{\phi}(y)|^2 |\Delta_{\underline{\mathbf{s}}}^*(y)|^{-1} dy < \infty.$$

Proof. Using (2.3), we have

$$\begin{split} C_{\phi} &= \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{P} |\langle \phi, \pi_{(u,v)} \phi \rangle_{L^{2}(V)}|^{2} d\mu_{l}(u,v) \\ &= \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{T} \left(\int_{V} |\phi * \tilde{\phi}_{t(v)}(u)|^{2} du \right) \prod_{j=1}^{r} v_{j}^{-d(r-1)-2} d\mu_{l}(t(v)) \\ &= \frac{1}{(2\pi)^{n}} \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{T} \left(\int_{\Omega_{\varepsilon}} |\hat{\phi}(y) \overline{\hat{\phi}(t(v)'y)}|^{2} dy \right) d\mu_{l}(t(v)) \\ &= \frac{1}{(2\pi)^{n}} \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{\Omega_{\varepsilon}} |\hat{\phi}(y)|^{2} \left(\int_{T} |\hat{\phi}(t(v)'y)|^{2} d\mu_{l}(t(v)) \right) dy. \end{split}$$

For $y \in \Omega_{\varepsilon}$, there exists $v^1 \in V_+$ such that $y = t(v^1)'P(\varepsilon)e$. Using (3.2) we obtain

$$\begin{split} &\int_{T} |\hat{\phi}(t(v)'y)|^{2} d\mu_{l}(t(v)) \\ &= \int_{T} \left| \hat{\phi}\Big((t(v^{1})t(v))'P(\varepsilon)e\Big) \right|^{2} d\mu_{l}(t(v)) \\ &= \int_{T} |\hat{\phi}(t(v)'P(\varepsilon)e)|^{2} d\mu_{l}(t(v)) \\ &= \int_{\Omega_{\varepsilon}} |\hat{\phi}(y)|^{2} |\Delta_{\underline{\mathbf{s}}}^{*}(y)|^{-1} dy. \end{split}$$

The proof of Theorem 1 is completed.

Suppose ϕ and ψ are admissible wavelets. We define the "inner product" of ϕ and ψ by

$$\langle \phi, \psi \rangle_{AW} = \int_{V} \hat{\phi}(y) \overline{\hat{\psi}(y)} |\Delta_{\underline{\mathbf{s}}}^{*}(y)|^{-1} dy.$$

Remark. If $\phi \in AW_{\varepsilon}$, $\psi \in AW_{\delta}$, $\varepsilon \neq \delta$, then $\langle \phi, \psi \rangle_{AW} = 0$. For $f \in H_{\varepsilon}$, $\phi \in AW_{\varepsilon}$, we define the wavelet transform of f with respect to ϕ by

$$W_{\phi}f(u,v) = \langle f, \pi_{(u,v)}\phi \rangle_{L^2(V)}.$$

Theorem 2. Suppose $f, g \in H_{\varepsilon}, \phi, \psi \in AW_{\varepsilon}$. Then

$$\langle W_{\phi}f, W_{\psi}g \rangle_{L^{2}(P, d\mu_{l})} = \langle \psi, \phi \rangle_{AW} \langle f, g \rangle_{L^{2}(V)}.$$

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In particular,

$$||W_{\phi}f||^{2}_{L^{2}(P,\,d\mu_{l})} = C_{\phi}||f||^{2}_{L^{2}(V)}.$$

Theorem 2 can be proved in a similar way as Theorem 1. From the theory of square-integrable representation of nonunimodular groups (cf [2]), Theorem 1 and Theorem 2 are equivalent.

We are going to construct a family of admissible wavelets which is complete and orthonormal with respect to $\langle \cdot, \cdot \rangle_{AW}$.

Let $\{c_{jk}^l: l = 1, \cdots, d\}$ be an orthonormal basis of V_{jk} . The set of indices \mathcal{A} is defined by

$$\mathcal{A} = \{ \alpha \in V : \ \alpha = \sum_{j=1}^{r} \alpha_j c_j + \sum_{j < k} \sum_{l=1}^{d} \alpha_{jk}^l c_{jk}^l, \ \alpha_j, \alpha_{jk}^l \text{ are nonnegative integers} \}.$$

Let $L_m^{(\mu)}(s)$ be the Laguerre polynomials defined by

$$L_m^{(\mu)}(s) = \sum_{j=0}^m \binom{m+\mu}{m-j} \frac{(-s)^j}{j!} = \frac{1}{m!} e^s s^{-\mu} \left(\frac{d}{ds}\right)^m (e^{-s} s^{m+\mu}), \quad \mu > -1,$$

and $H_m(s)$ be the Hermite polynomials defined by

$$H_m(s) = \sum_{j=0}^{\lfloor m/2 \rfloor} (-1)^j \frac{m!}{j!(m-2j)!} (2s)^{m-2j} = (-1)^m e^{s^2} \left(\frac{d}{ds}\right)^m (e^{-s^2}).$$

The Laguerre polynomials and the Hermite polynomials satisfy the following orthogonal relations respectively.

$$\int_{0}^{\infty} e^{-s} s^{\mu} L_{m}^{(\mu)}(s) L_{k}^{(\mu)}(s) \, ds = \Gamma(\mu+1) \binom{m+\mu}{m} \delta_{mk},$$
$$\int_{-\infty}^{\infty} e^{-s^{2}} H_{m}(s) H_{k}(s) \, ds = \pi^{\frac{1}{2}} 2^{m} m! \delta_{mk}.$$

And they are complete in $L^2(\mathbf{R}^+, e^{-s}s^{\mu}ds)$ and $L^2(\mathbf{R}, e^{-s^2}ds)$ respectively (cf [10], Chapter V). We also need following notations. For

$$\alpha = \sum_{j=1}^{r} \alpha_j c_j + \sum_{j < k} \sum_{l=1}^{d} \alpha_{jk}^l c_{jk}^l \in \mathcal{A},$$

we set

$$\begin{aligned} |\alpha| &= \sum_{j=1}^{r} \alpha_j + \sum_{j < k} \sum_{l=1}^{d} \alpha_{jk}^l, \\ \alpha! &= \prod_{j=1}^{r} \alpha_j! \prod_{j < k} \prod_{l=1}^{d} \alpha_{jk}^l!, \\ \alpha^0 &= (\alpha_1, \cdots, \alpha_r), \\ |\alpha^0| &= \sum_{j=1}^{r} \alpha_j, \\ \alpha^0! &= \prod_{j=1}^{r} \alpha_j!. \end{aligned}$$

For $\mathbf{s} \in \mathbf{C}^r$, $\lambda \in \mathbf{C}$, we write

 $\mathbf{s} + \lambda = (s_1 + \lambda, \cdots, s_r + \lambda).$

The gamma function of the symmetric cone Ω is defined by

$$\Gamma_{\Omega}(\mathbf{s}) = \int_{\Omega} e^{-\operatorname{tr}(x)} \Delta_{\mathbf{s}}(x) \Delta(x)^{-\frac{n}{r}} dx$$

= $(2\pi)^{\frac{n-r}{2}} \prod_{j=1}^{r} \Gamma(s_j - (j-1)\frac{d}{2}), \quad \operatorname{Re}s_j > (j-1)\frac{d}{2}, \ j = 1, \cdots, r.$

Assume that

$$v = \sum_{j=1}^{r} v_j c_j + \sum_{j < k} \sum_{l=1}^{d} v_{jk}^l c_{jk}^l \in V_+,$$

we define

$$\psi_{\alpha}(v) = C \prod_{j=1}^{r} e^{-v_{j}^{2}} v_{j}^{\nu - \frac{n}{r}} L_{\alpha_{j}}^{(\mu_{j})}(2v_{j}^{2}) \prod_{j < k} \prod_{l=1}^{d} e^{-\frac{1}{2}(v_{jk}^{l})^{2}} H_{\alpha_{jk}^{l}}(v_{jk}^{l})$$

where

$$C = 2^{\frac{1}{2}(\nu r - n + |\alpha| - |\alpha^{0}|)} \alpha^{0}! (\alpha!)^{-\frac{1}{2}} \Gamma_{\Omega} (\alpha^{0} + \nu - \frac{n}{r})^{-\frac{1}{2}},$$

$$\mu_{j} = \nu - \frac{n}{r} - 1 - \frac{d}{2}(j - 1), \quad \nu > 1 + d(r - 1).$$

We regard $\psi_{\alpha}(v)$ as the functions on group T. From the orthogonal relations and completeness of the Laguerre polynomials and Hermite polynomials, we conclude that $\{\psi_{\alpha}(v): \alpha \in \mathcal{A}\}$ is an orthonormal basis of $L^{2}(T, d\mu_{l})$.

Now we define a family of functions $\phi_{\alpha}^{\varepsilon}$ by

$$\hat{\phi}_{\alpha}^{\varepsilon}(y) = \begin{cases} \psi_{\alpha}(v), & y = t(v)' P(\varepsilon)e, \\ 0, & y \notin \Omega_{\varepsilon}. \end{cases}$$

Then $\phi_{\alpha}^{\varepsilon} \in H_{\varepsilon}$ is admissible and $\{\phi_{\alpha}^{\varepsilon} : \varepsilon \in E, \alpha \in \mathcal{A}\}$ is an complete orthonormal system with respect to $\langle \cdot, \cdot \rangle_{AW}$.

4. The decomposition of $L^2_{\nu}(T_{\Omega})$

The weighted L^2 -space on the symmetric tube domain T_{Ω} is defined by

$$L^{2}_{\nu}(T_{\Omega}) = \{F: \|F\|^{2}_{\nu} = \int_{T_{\Omega}} |F(x+iy)|^{2} \Delta(y)^{\nu - \frac{2n}{r}} \, dx \, dy < \infty\}.$$

 $\mathcal{H}^2_{\nu}(T_{\Omega})$, the weighted Bergman space, is the subspace of all holomorphic functions in $L^2_{\nu}(T_{\Omega})$, *i.e.*,

$$\mathcal{H}^2_{\nu}(T_{\Omega}) = \{F \in L^2_{\nu}(T_{\Omega}) : F \text{ is holomorphic on } T_{\Omega}\}.$$

We assume that $\nu > 1 + d(r-1)$ so that $\mathcal{H}^2_{\nu}(T_{\Omega}) \neq \{0\}$. We want to decompose $L^2_{\nu}(T_{\Omega})$ into the direct sum of subspaces such that the first component is exactly the weighted Bergman space $\mathcal{H}^2_{\nu}(T_{\Omega})$.

Suppose $f \in H_{\varepsilon}, \phi \in AW_{\varepsilon}$. For $\nu > 1 + d(r-1)$, we define the weighted wavelet transform W_{ϕ}^{ν} by

$$F(u + it(v)e) = W_{\phi}^{\nu} f(u + it(v)e) = C_{\phi}^{-\frac{1}{2}} W_{\phi} f(u, v) \Delta(t(v)e)^{-\frac{\nu}{2}}.$$

Set

$$\mathcal{H}^{\varepsilon}_{\alpha} = \{ F = W^{\nu}_{\phi^{\varepsilon}_{\alpha}} f : f \in H_{\varepsilon} \}.$$

Proposition 2. For $\nu > 1 + d(r-1)$, we have

(4.1)
$$L^2_{\nu}(T_{\Omega}) = \bigoplus_{\varepsilon \in E, \alpha \in \mathcal{A}} \mathcal{H}^{\varepsilon}_{\alpha}$$

and

$$\mathcal{H}_0^e = \mathcal{H}_\nu^2(T_\Omega).$$

Proof. By Theorem 2, W^{ν}_{ϕ} is an isometric operator from H_{ε} into $L^{2}_{\nu}(T_{\Omega})$ and $\mathcal{H}^{\varepsilon}_{\alpha}$'s are mutually orthogonal subspaces of $L^{2}_{\nu}(T_{\Omega})$. We need to prove

$$L^2_{\nu}(T_{\Omega}) \subset \bigoplus_{\varepsilon \in E, \alpha \in \mathcal{A}} \mathcal{H}^{\varepsilon}_{\alpha}.$$

Suppose $F \in L^2_{\nu}(T_{\Omega})$. Write $F_v(u) = F(u + it(v)e)$. For $v \in V_+$ almost every where, $F_v \in L^2(V)$. We let $G(v, y) = \hat{F}_v(y)$. Fix $y = t(v^1)'P(\varepsilon)e \in \Omega_{\varepsilon}$, then G(v, y), regarded as the function on T, is in $L^2(T, \Delta(t(v)e)^{\nu - \frac{n}{r}}d\mu_l(t(v)))$. Since

$$\{\Delta(t(v)e)^{-\frac{\nu}{2}+\frac{n}{2r}}\hat{\phi}^{\varepsilon}_{\alpha}(t(v)'y): \alpha \in \mathcal{A}\}$$

is an orthonormal basis of $L^2(T, \Delta(t(v)e)^{\nu-\frac{n}{r}}d\mu_l(t(v)))$, we get

$$G(v,y) = \sum_{\alpha \in \mathcal{A}} a_{\alpha}(y) \Delta(t(v)e)^{-\frac{\nu}{2} + \frac{n}{2r}} \hat{\phi}^{\varepsilon}_{\alpha}(t(v)'y), \quad y \in \Omega_{\varepsilon}$$

We define the functions f_{α}^{ε} by

$$\hat{f}^{\varepsilon}_{\alpha}(y) = \begin{cases} a_{\alpha}(y), & y \in \Omega_{\varepsilon}, \\ 0, & y \notin \Omega_{\varepsilon}. \end{cases}$$

It is easy to see that $f^{\varepsilon}_{\alpha} \in H_{\varepsilon}$. Therefore we have

$$F(u+it(v)e) = \sum_{\varepsilon \in E, \alpha \in \mathcal{A}} W^{\nu}_{\phi^{\varepsilon}_{\alpha}} f^{\varepsilon}_{\alpha}(u+it(v)e).$$

This proves (4.1).

Now we assume that

$$F(u+it(v)e) = W^{\nu}_{\phi^{\rho}_0} f(u+it(v)e), \quad f \in H_e.$$

Note that

$$\begin{split} \mathrm{tr}(y) &= \sum_{j=1}^r y_j = \sum_{j=1}^r v_j^2 + \frac{1}{2} \sum_{j < k} \sum_{l=1}^d |v_{jk}^l|^2, \\ \Delta(y) &= \prod_{j=1}^r v_j^2, \qquad y = t(v)'e, \end{split}$$

we have

$$\hat{\phi}_0^e(y) = \begin{cases} \Gamma_{\Omega}(\nu - \frac{n}{r})^{-\frac{1}{2}} 2^{\frac{\nu r - n}{2}} e^{-\operatorname{tr} y} \Delta(y)^{\frac{\nu}{2} - \frac{n}{2r}}, & y \in \Omega, \\ 0, & y \notin \Omega. \end{cases}$$

Therefore,

$$\begin{split} F(u+it(v)e) &= \Delta(t(v)e)^{-\frac{\nu}{2}-\frac{n}{2r}} \int_{V} f(x)\overline{\phi(t(v)^{-1}x-t(v)^{-1}u)} \, dx \\ &= (2\pi)^{-n} \Delta(t(v)e)^{-\frac{\nu}{2}+\frac{n}{2r}} \int_{\Omega} \hat{f}(y)e^{i(u|y)}\overline{\phi(t(v)'y)} \, dy \\ &= (2\pi)^{-n} \Gamma_{\Omega}(\nu-\frac{n}{r})^{-\frac{1}{2}} \int_{\Omega} e^{i(u+it(v)e|y)} \hat{f}(y)\Delta(2y)^{\frac{\nu}{2}-\frac{n}{2r}} \, dy, \end{split}$$

where we make use of the equality

$$\Delta(t(v)e) = \operatorname{Det}(t(v))^{\frac{r}{n}} = \operatorname{Det}(t(v)')^{\frac{r}{n}} = \Delta(t(v)'e).$$

We define the map \mathcal{F}_{ν} by

$$g(y) = \mathcal{F}_{\nu}f(y) = (2\pi)^{-\frac{n}{2}}\Gamma_{\Omega}(\nu - \frac{n}{r})^{-\frac{1}{2}}\hat{f}(y)\Delta(2y)^{-\frac{\nu}{2} + \frac{n}{2r}}$$

and the map \mathcal{L}_{ν} by

$$F(u+it(v)e) = \mathcal{L}_{\nu}g(u+it(v)e) = (2\pi)^{-\frac{n}{2}} \int_{\Omega} e^{i(u+it(v)e|y)}g(y)\Delta(2y)^{\nu-\frac{n}{r}} \, dy.$$

It is obvious that \mathcal{F}_{ν} is an isomorphism from $H^{2}(V)$ onto $L^{2}_{\nu}(\Omega) = L^{2}(\Omega, \Delta(2y)^{\nu-\frac{n}{r}}dy)$. \mathcal{L}_{ν} is an isomorphism from $L^{2}_{\nu}(\Omega)$ onto $\mathcal{H}^{2}_{\nu}(T_{\Omega})$ (cf [3]). We see that $W^{\nu}_{\phi^{e}_{0}} = \mathcal{L}_{\nu} \circ \mathcal{F}_{\nu}$ and

$$||F||_{L^{2}_{\nu}(T_{\Omega})}^{2} = \Gamma_{\Omega}(\nu - \frac{n}{r})||g||_{L^{2}_{\nu}(\Omega)}^{2} = ||f||_{L^{2}(V)}^{2}.$$

Clearly, \mathcal{H}_0^e is exactly the weighted Bergman space $\mathcal{H}_{\nu}^2(T_{\Omega})$.

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