# Wavelet Transforms and Symmetric Tube Domains 

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#### Abstract

We extend wavelet analysis to the symmetric tube domains and their Shilov boundaries. Our approach is based on the theory of Jordan algebra.


One-dimensional wavelet analysis can be explained in terms of square-integrable representation of the affine group (cf. [4], [6]). It is an intermediate between the function theory on the upper half-plane of one complex variable and the harmonic analysis on the real line (cf. [7], [9]). In this paper we extend wavelet analysis to the symmetric tube domains and their Shilov boundaries, the higher dimensional analogues of the upper half-plane and the real line. We assume that $V$ is a simple Euclidean Jordan algebra, $\Omega$ is the associated symmetric cone and $T_{\Omega}$ is the symmetric tube domain over $\Omega$. In $\S 1$, we recall some notations and facts about Jordan algebras and symmetric cones, especially the Iwasawa subgroup $P$ of the holomorphic automorphism group of $T_{\Omega} . P$ has a natural unitary representation $\pi$ on $L^{2}(V)$. In $\S 2$, we decompose $L^{2}(V)$ into the direct sum of the irreducible invariant closed subspaces under $\pi$. In $\S 3$, we give an explicit characterization of the admissibility condition in terms of Fourier transform and Jordan algebra. We also give a family of admissible wavelets, which is a complete orthonormal system in a sense. Finally in §4, we use wavelet transforms to decompose the weighted $L^{2}$-space on the tube domain $T_{\Omega}$ into a direct sum of subspaces such that the first component is exactly the weighted Bergman space.

A good reference on Jordan algebras, symmetric cones and tube domains is the book [3] by J. Faraut and A. Korányi. Various authors developed the theory of continuous wavelet in view of square-integrable group representations, for example, in [5], [8] and in particular [1].

## 1. Iwasawa subgroup

Throughout this paper we keep the following assumptions and notations, which are the same as in [3].
$V$ is an $n$-dimensional simple Euclidean Jordan algebra with identity $e$. $x y$ denotes the Jordan product of $x$ and $y . \operatorname{tr}(x)$ and $\operatorname{det}(x)$ are defined as in [3]. We also write $\Delta(x)$ instead of $\operatorname{det}(x)$. The inner product on $V$ is given by $(x \mid y)=\operatorname{tr}(x y) . L(x)$ is the linear map of $V$ defined by $L(x) y=x y$. An element $c \in V$ is idempotent if $c^{2}=c$. The only eigenvalues of $L(c)$ are $1, \frac{1}{2}$, and 0 . The corresponding eigenspaces are denoted by $V(c, 1), V\left(c, \frac{1}{2}\right)$ and $V(c, 0)$. We fix a Jordan frame $\left\{c_{1}, \cdots, c_{r}\right\}$, where $r$ is the rank of $V$. Then we have the Peirce decomposition

$$
V=\bigoplus_{j \leq k} V_{j k}
$$

where

$$
\begin{aligned}
& V_{i i}=V\left(c_{i}, 1\right)=\mathbf{R} c_{i}, \\
& V_{i j}=V\left(c_{i}, \frac{1}{2}\right) \cap V\left(c_{j}, \frac{1}{2}\right) .
\end{aligned}
$$

$d=\operatorname{dim} V_{i j}$, which does not depend on $i$ and $j$, is called the degree of $V$. Let

$$
P(x)=2 L(x)^{2}-L\left(x^{2}\right)
$$

be the quadratic representation, and write

$$
x \square y=L(x y)-[L(x), L(y)] .
$$

For given $j$ and for $z^{(j)} \in \bigoplus_{k=j+1}^{r} V_{j k}$ the Frobenius transform $\tau\left(z^{(j)}\right)$ is defined by

$$
\tau\left(z^{(j)}\right)=\exp \left(2 z^{(j)} \square c_{j}\right)
$$

Let $\Omega$ be the symmetric cone which consists of elements $x$ in $V$ such that $L(x)$ is positive definite. $G(\Omega)$ denotes the automorphism group of $\Omega$ and $G$ is the identity component of $G(\Omega) . G$ has Iwasawa decomposition $G=N A K$, where

$$
\begin{aligned}
& K=\{g \in G: g e=e\}, \\
& A=\left\{P(a): a=\sum_{j=1}^{r} a_{j} c_{j}, a_{j}>0\right\}, \\
& N=\left\{\tau\left(z^{(1)}\right) \cdots \tau\left(z^{(r-1)}\right): z^{(j)} \in \bigoplus_{k=j+1}^{r} V_{j k}\right\}
\end{aligned}
$$

are compact, diagonal and strict triangular respectively. $A$ normalizes $N$ and

$$
\begin{equation*}
P(a) \tau\left(z^{(j)}\right)=\tau\left(\tilde{z}^{(j)}\right) P(a) \tag{1.1}
\end{equation*}
$$

where

$$
\begin{aligned}
& z^{(j)}=\sum_{j<k} z_{j k}, \quad z_{j k} \in V_{j k}, \\
& \tilde{z}^{(j)}=\sum_{j<k} \tilde{z}_{j k}, \quad \tilde{z}_{j k} \in V_{j k}, \\
& \tilde{z}_{j k}=\frac{a_{k}}{a_{j}} z_{j k} .
\end{aligned}
$$

$T=N A$ is a semi-direct product. We will use another parametrization of the triangular subgroup $T$. Set

$$
V_{+}=\left\{u=\sum_{j=1}^{r} u_{j} c_{j}+\sum_{j<k} u_{j k}: u_{j}>0, u_{j k} \in V_{j k}\right\} .
$$

For $u \in V_{+}$, we define

$$
t(u)=P\left(b_{1}\right) \tau\left(u^{(1)}\right) P\left(b_{2}\right) \cdots \tau\left(u^{(r-1)}\right) P\left(b_{r}\right),
$$

where

$$
\begin{aligned}
& b_{j}=c_{1}+\cdots+c_{j-1}+u_{j} c_{j}+c_{j+1}+\cdots+c_{r}, \\
& u^{(j)}=\sum_{k=j+1}^{r} u_{j k} .
\end{aligned}
$$

Then

$$
T=\left\{t(u): u \in V_{+}\right\} .
$$

Using (1.1), it is easy to determine the left and right Haar measures of $T$. The left Haar measure of $T$ is given by

$$
d \mu_{l}(t(u))=2^{r} \prod_{j=1}^{r} u_{j}^{-d(j-1)-1} d u
$$

and the right Haar measure of $T$ is given by

$$
d \mu_{r}(t(u))=2^{r} \prod_{j=1}^{r} u_{j}^{-d(r-j)-1} d u
$$

$T$ acts simply and transitively on $\Omega$. If

$$
x=\sum_{j=1}^{r} x_{j} c_{j}+\sum_{j<k} x_{j k}
$$

is the Peirce decomposition of $x=t(u) e$, then

$$
\begin{aligned}
& x_{j}=u_{j}^{2}+\frac{1}{2} \sum_{k=1}^{j-1}\left\|u_{k j}\right\|^{2}, \\
& x_{j k}=u_{j} u_{j k}+2 \sum_{l=1}^{j-1} u_{l j} u_{l k} .
\end{aligned}
$$

We identify $\Omega$ with $T$ by identification of $x=t(u) e$ and $t(u)$. Then we have

$$
\begin{aligned}
& d x=2^{r} \prod_{j=1}^{r} u_{j}^{d(r-j)+1} d u, \\
& \Delta(x)=\prod_{j=1}^{r} u_{j}^{2} .
\end{aligned}
$$

Therefore

$$
\Delta(x)^{-\frac{n}{r}} d x=d \mu_{l}(t(u))
$$

which gives the $G$-invariant measure on $\Omega$.
Let $V^{C}$ denote the complexification of $V . T_{\Omega}=V+i \Omega$ is the tube domain over $\Omega$ in $V^{C} . G\left(T_{\Omega}\right)$ denotes the holomorphic automorphism group of $T_{\Omega}$ and $G\left(T_{\Omega}\right)^{0}$ is the identity component of $G\left(T_{\Omega}\right)$. The Iwasawa decomposition of $G\left(T_{\Omega}\right)^{0}$ is given by $G\left(T_{\Omega}\right)^{0}=\underline{N} A \underline{K}$, where

$$
\begin{aligned}
& \underline{K}=\left\{g \in G\left(T_{\Omega}\right)^{0}: g(i e)=i e\right\} \supset K, \\
& \underline{N}=N^{+} N, \\
& N^{+}=\left\{\tau_{u}: z \mapsto z+u, \quad u \in V\right\} \cong V .
\end{aligned}
$$

Therefore,

$$
G\left(T_{\Omega}\right)^{0}=N^{+} T \underline{K} .
$$

We call it the partial Iwasawa decomposition as in Terras' book [11]. $T$ normalizes $N^{+}$as

$$
t(v) \tau_{u}=\tau_{t(v) u} t(v), \quad u \in V, v \in V_{+} .
$$

$P=\underline{N} A=N^{+} T$ is called the Iwasawa subgroup. $P$ is a nonunimodular group. Using the parametrization $(u, v)$ for $\tau_{u} t(v) \in P$, the left Haar measure of $P$ is given by

$$
d \mu_{l}(u, v)=2^{r} \prod_{j=1}^{r} v_{j}^{-d(r+j-2)-3} d u d v=\prod_{j=1}^{r} v_{j}^{-d(r-1)-2} d u d \mu_{l}(t(v))
$$

and the right Haar measure of $P$ is given by

$$
d \mu_{r}(u, v)=2^{r} \prod_{j=1}^{r} v_{j}^{-d(r-j)-1} d u d v=d u d \mu_{r}(t(v))
$$

$P$ acts on $T_{\Omega}$ simply and transitively. We identify $T_{\Omega}$ with $P$ by identification of $\tau_{u} t(v)(i e)$ and $\tau_{u} t(v)$. If $x+i y=\tau_{u} t(v)(i e)=u+i t(v) e$, then

$$
\Delta(y)^{-\frac{2 n}{r}} d x d y=d \mu_{l}(u, v)
$$

which is the $G\left(T_{\Omega}\right)^{0}$-invariant measure on $T_{\Omega}$. Note that

$$
\operatorname{Det}(g)=\Delta(g e)^{\frac{n}{r}}, \quad g \in G .
$$

$P$ has a natural unitary representation on $L^{2}(V)$ defined by

$$
\pi_{(u, v)}: f(x) \mapsto \Delta(t(v) e)^{-\frac{n}{2 r}} f\left(t(v)^{-1} x-t(v)^{-1} u\right)
$$

We shall decompose $L^{2}(V)$ into the direct sum of irreducible invariant closed subspaces under $\pi$.

## 2. The decomposition of $L^{2}(V)$

In order to decompose $L^{2}(V)$, we need to identify the non-degenerate $T$-orbits of $V$ under the contragredient action of $T$, which is given by $x \mapsto t(v)^{\prime-1} x$ where $t(v)^{\prime}$ denotes the transpose of $t(v)$. First we prove

Lemma 1. (1) Suppose $z_{i j} \in V_{i j}, w_{k l} \in V_{k l}, i<j, k<l, i \neq l, k \neq j$, then

$$
\left[z_{i j} \square c_{i}, w_{k l} \square c_{k}\right]=0 .
$$

(2) Suppose $z_{i j} \in V_{i j}$, then

$$
\left(z_{i j} \square c_{i}\right)^{\prime}=z_{i j} \square c_{j} .
$$

Proof. (a) To prove (1), we use the facts

$$
\begin{aligned}
& V_{i j} \cdot V_{j k} \subset V_{i k}, \quad \text { if } \quad i \neq k, \\
& V_{i j} \cdot V_{k l}=\{0\}, \quad \text { if } \quad\{i, j\} \cap\{k, l\}=\emptyset, \\
& x y=\frac{1}{2}(x \mid y)\left(c_{i}+c_{j}\right), \quad \text { if } \quad x, y \in V_{i j}
\end{aligned}
$$

( cf [3], Theorem IV.2.1 (iii) and Proposition IV.1.4 (i) ). We also use the matrix of $z \square c$ with respect to the Peirce decomposition, when $c$ is idempotent in $V$ and $z \in V\left(c, \frac{1}{2}\right)$ ( see [3], proof of Lemma VI.3.1 ). Let

$$
x=\sum_{j=1}^{r} x_{j} c_{j}+\sum_{j<k} x_{j k}, \quad x_{j k} \in V_{j k}
$$

We compute separately in four cases.

1) If $k=i, l=j, i<j$, then

$$
\left(z_{i j} \square c_{i}\right)\left(w_{i j} \square c_{i}\right) x=\frac{x_{i}}{4}\left(z_{i j} \mid w_{i j}\right) c_{j}=\left(w_{i j} \square c_{i}\right)\left(z_{i j} \square c_{i}\right) x .
$$

2) If $k=i, l \neq j, i<j, l$, then

$$
\left(z_{i j} \square c_{i}\right)\left(w_{i l} \square c_{i}\right) x=\frac{x_{i}}{2} z_{i j} w_{i l}=\left(w_{i l} \square c_{i}\right)\left(z_{i j} \square c_{i}\right) x .
$$

3) If $k \neq i, l=j, i, k<j$, then

$$
\left(z_{i j} \square c_{i}\right)\left(w_{k j} \square c_{k}\right) x=\frac{1}{2}\left(z_{i j} \mid x_{i k} w_{k j}\right) c_{j}=\frac{1}{2}\left(w_{k j} \mid x_{i k} z_{i j}\right) c_{j}=\left(w_{k j} \square c_{k}\right)\left(z_{i j} \square c_{i}\right) x,
$$

where the second equality is due to the associativity of the inner product.
4) If $k \neq i, j, l \neq i, j, i<j, k<l$, we may assume $i<k$, then

$$
\left(z_{i j} \square c_{i}\right)\left(w_{k l} \square c_{k}\right) x=z_{i j}\left(x_{i k} w_{k l}\right)=w_{k l}\left(x_{i k} z_{i j}\right)=\left(w_{k l} \square c_{k}\right)\left(z_{i j} \square c_{i}\right) x
$$

where the second equality follows from the Lemma V.3.2 in [3].
(b) Take $x=z_{i j}, y=c_{i}+c_{j}$ in the identity

$$
\left[L(x), L\left(y^{2}\right)\right]+2[L(y), L(x y)]=0
$$

( cf [3]. Proposition II.1.1 ), we obtain

$$
\left[L\left(c_{i}\right), L\left(z_{i j}\right)\right]=\left[L\left(z_{i j}\right), L\left(c_{j}\right)\right]
$$

It follows that

$$
\left(z_{i j} \square c_{i}\right)^{\prime}=c_{i} \square z_{i j}=z_{i j} \square c_{j} .
$$

Let $z_{j k} \in V_{j k}(j<k)$ and put

$$
z^{(j)}=\sum_{k=j+1}^{r} z_{j k}, \quad z_{(k)}=\sum_{j=1}^{k-1} z_{j k}
$$

Put

$$
\tau^{\prime}\left(z_{(k)}\right)=\exp \left(2 z_{(k)} \square c_{k}\right) .
$$

If $z_{i j} \in V_{i j}, w_{k l} \in V_{k l}, i<j, k<l, i \neq l, k \neq j$, Lemma 1 implies that

$$
\tau\left(z_{i j}\right) \tau\left(w_{k l}\right)=\tau\left(w_{k l}\right) \tau\left(z_{i j}\right)
$$

and

$$
\tau\left(z_{i j}\right)^{\prime}=\tau^{\prime}\left(z_{i j}\right)
$$

Thus $\tau^{\prime}\left(z_{i j}\right)$ is a dual Frobenius transform. Also, by Lemma 1,

$$
\begin{aligned}
\tau\left(z^{(j)}\right) & =\tau\left(z_{j, j+1}\right) \cdots \tau\left(z_{j, r}\right), \\
\tau^{\prime}\left(z_{(k)}\right) & =\tau^{\prime}\left(z_{1, k}\right) \cdots \tau^{\prime}\left(z_{k-1, k}\right) .
\end{aligned}
$$

Therefore, for

$$
u=\sum_{j=1}^{r} u_{j} c_{j}+\sum_{j<k} u_{j k}, \quad u_{j}>0, u_{j k} \in V_{j k}
$$

we have, by also using (1.1),

$$
\begin{aligned}
t(u) & =P\left(b_{1}\right) \tau\left(u^{(1)}\right) P\left(b_{2}\right) \cdots \tau\left(u^{(r-1)}\right) P\left(b_{r}\right) \\
& =P\left(b_{1}\right) \tau\left(u_{12}\right) P\left(b_{2}\right) \tau\left(u_{13}\right) \tau\left(u_{23}\right) \cdots P\left(b_{r-1}\right) \tau\left(u_{1 r}\right) \cdots \tau\left(u_{r-1, r}\right) P\left(b_{r}\right) . \\
t(u)^{\prime} & =P\left(b_{r}\right) \tau^{\prime}\left(u_{r-1, r}\right) \tau^{\prime}\left(u_{r-2, r}\right) \cdots \tau^{\prime}\left(u_{1 r}\right) P\left(b_{r-1}\right) \cdots P\left(b_{2}\right) \tau^{\prime}\left(u_{12}\right) P\left(b_{1}\right) \\
& =P\left(b_{r}\right) \tau^{\prime}\left(u_{(r)}\right) P\left(b_{r-1}\right) \cdots \tau^{\prime}\left(u_{(2)}\right) P\left(b_{1}\right)
\end{aligned}
$$

where

$$
u_{(k)}=\sum_{j=1}^{k-1} u_{j k} .
$$

For $j=1, \cdots, r$, let $V^{(j)}$ be the subalgebra $V\left(c_{1}+\cdots+c_{j}, 1\right)$ of $V$ and $W^{(j)}$ be the subalgebra $V\left(c_{r-j+1}+\cdots+c_{r}, 1\right)$ of $V . P_{j}$ and $P_{j}^{*}$ denote the orthogonal projections onto $V^{(j)}$ and $W^{(j)}$ respectively. $\operatorname{det}_{(j)}$ and $\operatorname{det}_{(j)}^{*}$ are the determinants relative to $V^{(j)}$ and $W^{(j)}$ respectively. We define

$$
\begin{aligned}
\Delta_{j}(x) & =\operatorname{det}_{(j)}\left(P_{j} x\right), \\
\Delta_{j}^{*}(x) & =\operatorname{det}_{(j)}^{*}\left(P_{j}^{*} x\right) .
\end{aligned}
$$

Furthermore, for $\mathbf{s}=\left(s_{1}, \cdots, s_{r}\right)$. We let

$$
\begin{aligned}
& \Delta_{\mathbf{s}}(x)=\Delta_{1}(x)^{s_{1}-s_{2}} \cdots \Delta_{r-1}(x)^{s_{r-1}-s_{r}} \Delta_{r}(x)^{s_{r}} \\
& \Delta_{\mathbf{s}}^{*}(x)=\Delta_{1}^{*}(x)^{s_{1}-s_{2}} \cdots \Delta_{r-1}^{*}(x)^{s_{r-1}-s_{r}} \Delta_{r}^{*}(x)^{s_{r}} .
\end{aligned}
$$

For $x \in V, t(u) \in T$, we have

$$
\begin{equation*}
\Delta_{\mathbf{s}}^{*}\left(t(u)^{\prime} x\right)=u_{1}^{2 s_{r}} \cdots u_{r}^{2 s_{1}} \Delta_{\mathbf{s}}^{*}(x)=\Delta_{\mathbf{s}}^{*}\left(t(u)^{\prime} e\right) \Delta_{\mathbf{s}}^{*}(x) \tag{2.1}
\end{equation*}
$$

In particular, $\Delta_{\mathrm{s}}^{*}$ is invariant under the Frobenius transform $\tau^{\prime}\left(z_{(k)}\right)(\operatorname{cf}[3]$, Proposition VII.1.5 ).

Set

$$
\begin{aligned}
& E=\left\{\varepsilon=\sum_{j=1}^{r} \varepsilon_{j} c_{j}: \varepsilon_{j}=1 \text { or } i\right\} \\
& \Omega_{\varepsilon}=\left\{x \in V: x=t(u)^{\prime} P(\varepsilon) e, u \in V_{+}\right\} .
\end{aligned}
$$

Lemma 2. (1) The $\Omega_{\varepsilon}$ 's are disjoint and simply transitive orbits under the contragredient action of $T$. (2) $\bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$ is a set with a complementary of measure zero.

Proof. (a) Suppose that

$$
t(u)^{\prime} P(\varepsilon) e=t(v)^{\prime} P(\delta) e, \quad u, v \in V_{+}, \varepsilon, \delta \in E
$$

Write

$$
g=P(\delta) t(v)^{\prime-1} t(u)^{\prime} P(\varepsilon) .
$$

Since $t(u), t(v)$ are triangular and $P(\varepsilon), P(\delta)$ are diagonal, $g$ is triangular. On the other hand, since $g e=e$, from the Proposition VIII.2.4 in [3] $g$ is an automorphism of $V^{C}$ and $g^{\prime}=g^{-1}$. Therefore $g$ is diagonal. Because $t(u), t(v)$ have positive diagonal elements and $P(\varepsilon), P(\delta)$ have diagonal elements $1,-1$ or $i$, it is concluded that $u=v, \varepsilon=\delta$.
(b) Set

$$
B=\left\{x \in V: \Delta_{k}^{*}(x) \neq 0, k=1, \cdots, r\right\} .
$$

Obviously, $V \backslash B$ is a zero measure set. We will prove that $B=\bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$. It is easy to see that $B \supset \bigcup_{\varepsilon \in E} \Omega_{\varepsilon}$. Assume that

$$
x=\sum_{j=1}^{r} x_{j} c_{j}+\sum_{j<k} x_{j k} \in B .
$$

By [3], Theorem VI.3.5 we can write

$$
x=\tau^{\prime}\left(z_{(r)}\right) \cdots \tau^{\prime}\left(z_{(2)}\right) \sum_{j=1}^{r} a_{j} c_{j}
$$

where

$$
\begin{aligned}
& z_{(k)}=\sum_{j=1}^{k-1} z_{j k} \in \bigoplus_{j=1}^{k-1} V_{j k}, \\
& a_{j}=\frac{\Delta_{r-j+1}^{*}(x)}{\Delta_{r-j}^{*}(x)} \neq 0, \quad j=1, \cdots, r-1, \\
& a_{r}=\Delta_{1}^{*}(x) \neq 0 .
\end{aligned}
$$

Set

$$
\begin{aligned}
& \varepsilon_{j}=\left\{\begin{array}{lll}
1, & \text { if } & a_{j}>0, \\
i, & \text { if } & a_{j}<0,
\end{array}\right. \\
& u_{j}=\sqrt{\left|a_{j}\right|}, \\
& u_{j k}=u_{k} z_{j k} .
\end{aligned}
$$

Then, by (1.1),

$$
x=t(u)^{\prime} P(\varepsilon) e
$$

Remark. Clearly, $\Omega_{e}=\Omega, \Omega_{i e}=-\Omega . \Omega_{\varepsilon}$ is a connected open set in $V$ because $\Omega_{\varepsilon}$ is homeomorphic to $V_{+}$. But $\Omega_{\varepsilon}$ may not be convex neither $K$ invariant in general.

A simple example of Lemma 2 can be given as follows. Let $V$ be the space $\operatorname{Sym}(m, \mathbf{R})$ of all $m \times m$ symmetric matrices and $c_{j}=\operatorname{diag}(0, \cdots, 0,1,0, \cdots, 0)$. An element $t$ in $T$ has the following form: $t x=u x u^{\prime}$, where $u$ is a lower triangular matrix with positive diagonal elements. Let $\Sigma$ denote the set of all diagonal matrices with diagonal elements $\pm 1 . \Omega_{\sigma}(\sigma \in \Sigma)$ consists of all matrices of form $u^{\prime} \sigma u$. Then $\Omega_{\sigma}$ 's are disjoint and simply transitive orbits under the adjoint action of $T$ and $\bigcup_{\sigma \in \Sigma} \Omega_{\sigma}$ is a total measure set. Now we are ready to decompose $L^{2}(V)$. Set

$$
H_{\varepsilon}=\left\{f \in L^{2}(V): \operatorname{supp} \hat{f} \subseteq \mathrm{Cl}\left(\Omega_{\varepsilon}\right)\right\} .
$$

Proposition 1. Each of $H_{\varepsilon}$ is an irreducible invariant closed subspace of $L^{2}(V)$ under $\pi$ and

$$
\begin{equation*}
L^{2}(V)=\bigoplus_{\varepsilon \in E} H_{\varepsilon} . \tag{2.2}
\end{equation*}
$$

Proof. (2.2) follows from Lemma 2. Because

$$
\left(\pi_{(u, v)} f\right)^{\wedge}(y)=\Delta(t(v) e)^{\frac{n}{2 r}} e^{-i(u \mid y)} \hat{f}\left(t(v)^{\prime} y\right),
$$

it is easy to see that $H_{\varepsilon}$ is invariant under $\pi$. We need to prove that $H_{\varepsilon}$ is irreducible. Let $W$ be a non-zero invariant closed subspace of $H_{\varepsilon}$ under $\pi$ and $W^{+}$the orthogonal complement of $W$ in $H_{\varepsilon}$. Taking a function $g \in W$, not identically zero, if $f \in W^{+}$, then

$$
\left\langle f, \pi_{(u, v)} g\right\rangle_{L^{2}(V)}=\int_{V} f(x) \overline{\pi_{(u, v)} g(x)} d x=0, \quad u \in V, v \in V_{+} .
$$

Write

$$
\begin{aligned}
& \tilde{g}(x)=\overline{g(-x)} \\
& g_{t(v)}(x)=\Delta(t(v) e)^{-\frac{n}{2 r}} g\left(t(v)^{-1} x\right)
\end{aligned}
$$

We have

$$
\begin{equation*}
\left\langle f, \pi_{(u, v)} g\right\rangle_{L^{2}(V)}=f * \tilde{g}_{t(v)}(u) . \tag{2.3}
\end{equation*}
$$

Therefore,

$$
\begin{equation*}
\left(f * \tilde{g}_{t(v)}\right)^{\wedge}(y)=\Delta(t(v) e)^{\frac{n}{2 r}} \hat{f}(y) \overline{\hat{g}\left(t(v)^{\prime} y\right)}=0, \quad \text { a.e. } y \in V . \tag{2.4}
\end{equation*}
$$

Set

$$
\begin{aligned}
& S_{1}=\operatorname{supp} \hat{f} \cap \Omega_{\varepsilon}, \\
& S_{2}=\operatorname{supp} \hat{g} \cap \Omega_{\varepsilon} .
\end{aligned}
$$

$S_{1}^{d}$ and $S_{2}^{d}$ consist of points of density of $S_{1}$ and $S_{2}$ respectively. $S_{2}^{d}$ is a positive measure set since $g$ is not identically zero. If $S_{1}^{d}$ has positive measure, by Lemma 2 , there exists $t\left(v_{0}\right) \in T$ such that $S=S_{1}^{d} \cap t\left(v_{0}\right)^{\prime-1} S_{2}^{d}$ has positive measure. But

$$
\left(f * \tilde{g}_{t(v)}\right)^{\wedge}(y) \neq 0, \quad y \in S,
$$

which contradicts (2.4). Therefore $f$ is identically zero. This proves that $H_{\varepsilon}$ is irreducible.

Remark. For $F$ in $H^{2}\left(T_{\Omega}\right)$, the Hardy space on $T_{\Omega}$, the following limit exists,

$$
\lim _{y \rightarrow 0, y \in \Omega} F(\cdot+i y)=f, \quad \text { in } L^{2}(V) .
$$

Then

$$
H^{2}(V)=\left\{f \in L^{2}(V): \text { there exists } F \in H^{2}\left(T_{\Omega}\right) \text { such that } f=\lim F\right\}
$$

is called the Hardy space on $V$. It is easy to see that $H_{e}=H^{2}(V)$ and $H_{i e}=\overline{H^{2}(V)}$.

## 3. The admissibility condition

The restriction of $\pi$ on $H_{\varepsilon}$ is square-integrable, i.e., there exists a function $\phi(\neq 0)$ in $H_{\varepsilon}$ such that

$$
\begin{equation*}
C_{\phi}=\frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{P}\left|\left\langle\phi, \pi_{(u, v)} \phi\right\rangle_{L^{2}(V)}\right|^{2} d \mu_{l}(u, v)<\infty \tag{3.1}
\end{equation*}
$$

(3.1) is called the admissibility condition and $\phi$ is called an admissible wavelet. We want to give a characterization of the admissibility conditioin in terms of Fourier transform and Jordan algebra, which does not involve any group representation.

Lemma 3. Suppose $x=t(u)^{\prime} P(\varepsilon) e$ in $\Omega_{\varepsilon}$. If

$$
x=\sum_{j=1}^{r} x_{j} c_{j}+\sum_{j<k} x_{j k}
$$

is the Peirce decomposition of $x$, then

$$
\begin{aligned}
& x_{j}=\varepsilon_{j}^{2} u_{j}^{2}+\frac{1}{2} \sum_{k=j+1}^{r} \varepsilon_{k}^{2}\left\|u_{j k}\right\|^{2}, \\
& x_{j k}=\varepsilon_{k}^{2} u_{k} u_{j k}+2 \sum_{l=k+1}^{r} \varepsilon_{l}^{2} u_{j l} u_{k l} .
\end{aligned}
$$

Lemma 3 can be proved in a similar way as in [3], Proposition VI.3.8.
For the transformation $x=t(u)^{\prime} P(\varepsilon) e$, by Lemma 3, it is easy to compute that

$$
\begin{aligned}
d x & =2^{r} \prod_{j=1}^{r} u_{j}^{d(j-1)+1} d u \\
& =\prod_{j=1}^{r} u_{j}^{2 d(j-1)+2} d \mu_{l}(t(u))
\end{aligned}
$$

Let

$$
\underline{\mathbf{s}}=(1+d(r-1), 1+d(r-2), \cdots, 1) .
$$

By (2.1),

$$
\Delta_{\underline{\mathbf{s}}}^{*}(x)=\Delta_{\underline{\mathbf{s}}}^{*}\left(t(u)^{\prime} e\right) \Delta_{\underline{\mathbf{s}}}^{*}(P(\varepsilon) e)
$$

Therefore,

$$
\left|\Delta_{\underline{\mathbf{s}}}^{*}(x)\right|=\Delta_{\underline{\mathbf{s}}}^{*}\left(t(u)^{\prime} e\right) .
$$

and we have

$$
\begin{equation*}
\left|\Delta_{\underline{\mathbf{s}}}^{*}(x)\right|^{-1} d x=d \mu_{l}(t(u)) . \tag{3.2}
\end{equation*}
$$

We denoted by $A W_{\varepsilon}$ the set of all admissible wavelets in $H_{\varepsilon}$.

Theorem 1. Suppose $\phi(\neq 0)$ in $H_{\varepsilon}$. Then $\phi \in A W_{\varepsilon}$ if and only if

$$
C_{\phi}=\int_{\Omega_{\varepsilon}}|\hat{\phi}(y)|^{2}\left|\Delta_{\underline{\mathbf{s}}}^{*}(y)\right|^{-1} d y<\infty .
$$

Proof. Using (2.3), we have

$$
\begin{aligned}
C_{\phi} & =\frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{P}\left|\left\langle\phi, \pi_{(u, v)} \phi\right\rangle_{L^{2}(V)}\right|^{2} d \mu_{l}(u, v) \\
& =\frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{T}\left(\int_{V}\left|\phi * \tilde{\phi}_{t(v)}(u)\right|^{2} d u\right) \prod_{j=1}^{r} v_{j}^{-d(r-1)-2} d \mu_{l}(t(v)) \\
& =\frac{1}{(2 \pi)^{n}} \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{T}\left(\int_{\Omega_{\varepsilon}}\left|\hat{\phi}(y) \overline{\hat{\phi}\left(t(v)^{\prime} y\right)}\right|^{2} d y\right) d \mu_{l}(t(v)) \\
& =\frac{1}{(2 \pi)^{n}} \frac{1}{\|\phi\|_{L^{2}(V)}^{2}} \int_{\Omega_{\varepsilon}}|\hat{\phi}(y)|^{2}\left(\int_{T}\left|\hat{\phi}\left(t(v)^{\prime} y\right)\right|^{2} d \mu_{l}(t(v))\right) d y
\end{aligned}
$$

For $y \in \Omega_{\varepsilon}$, there exists $v^{1} \in V_{+}$such that $y=t\left(v^{1}\right)^{\prime} P(\varepsilon) e$. Using (3.2) we obtain

$$
\begin{aligned}
& \int_{T}\left|\hat{\phi}\left(t(v)^{\prime} y\right)\right|^{2} d \mu_{l}(t(v)) \\
= & \int_{T}\left|\hat{\phi}\left(\left(t\left(v^{1}\right) t(v)\right)^{\prime} P(\varepsilon) e\right)\right|^{2} d \mu_{l}(t(v)) \\
= & \int_{T}\left|\hat{\phi}\left(t(v)^{\prime} P(\varepsilon) e\right)\right|^{2} d \mu_{l}(t(v)) \\
= & \int_{\Omega_{\varepsilon}}|\hat{\phi}(y)|^{2}\left|\Delta_{\underline{\mathbf{s}}}^{*}(y)\right|^{-1} d y .
\end{aligned}
$$

The proof of Theorem 1 is completed.
Suppose $\phi$ and $\psi$ are admissible wavelets. We define the "inner product" of $\phi$ and $\psi$ by

$$
\langle\phi, \psi\rangle_{A W}=\int_{V} \hat{\phi}(y) \overline{\hat{\psi}(y)}\left|\Delta_{\underline{\mathbf{s}}}^{*}(y)\right|^{-1} d y .
$$

Remark. If $\phi \in A W_{\varepsilon}, \psi \in A W_{\delta}, \varepsilon \neq \delta$, then $\langle\phi, \psi\rangle_{A W}=0$. For $f \in H_{\varepsilon}, \phi \in$ $A W_{\varepsilon}$, we define the wavelet transform of $f$ with respect to $\phi$ by

$$
W_{\phi} f(u, v)=\left\langle f, \pi_{(u, v)} \phi\right\rangle_{L^{2}(V)} .
$$

Theorem 2. Suppose $f, g \in H_{\varepsilon}, \phi, \psi \in A W_{\varepsilon}$. Then

$$
\left\langle W_{\phi} f, W_{\psi} g\right\rangle_{L^{2}\left(P, d \mu_{l}\right)}=\langle\psi, \phi\rangle_{A W}\langle f, g\rangle_{L^{2}(V)} .
$$

In particular,

$$
\left\|W_{\phi} f\right\|_{L^{2}\left(P, d \mu_{l}\right)}^{2}=C_{\phi}\|f\|_{L^{2}(V)}^{2}
$$

Theorem 2 can be proved in a similar way as Theorem 1. From the theory of square-integrable representation of nonunimodular groups (cf [2] ), Theorem 1 and Theorem 2 are equivalent.

We are going to construct a family of admissible wavelets which is complete and orthonormal with respect to $\langle\cdot, \cdot\rangle_{A W}$.

Let $\left\{c_{j k}^{l}: l=1, \cdots, d\right\}$ be an orthonormal basis of $V_{j k}$. The set of indices $\mathcal{A}$ is defined by
$\mathcal{A}=\left\{\alpha \in V: \alpha=\sum_{j=1}^{r} \alpha_{j} c_{j}+\sum_{j<k} \sum_{l=1}^{d} \alpha_{j k}^{l} c_{j k}^{l}, \alpha_{j}, \alpha_{j k}^{l}\right.$ are nonnegative integers $\}$.
Let $L_{m}^{(\mu)}(s)$ be the Laguerre polynomials defined by

$$
L_{m}^{(\mu)}(s)=\sum_{j=0}^{m}\binom{m+\mu}{m-j} \frac{(-s)^{j}}{j!}=\frac{1}{m!} e^{s} s^{-\mu}\left(\frac{d}{d s}\right)^{m}\left(e^{-s} s^{m+\mu}\right), \quad \mu>-1,
$$

and $H_{m}(s)$ be the Hermite polynomials defined by

$$
H_{m}(s)=\sum_{j=0}^{[m / 2]}(-1)^{j} \frac{m!}{j!(m-2 j)!}(2 s)^{m-2 j}=(-1)^{m} e^{s^{2}}\left(\frac{d}{d s}\right)^{m}\left(e^{-s^{2}}\right)
$$

The Laguerre polynomials and the Hermite polynomials satisfy the following orthogonal relations respectively.

$$
\begin{aligned}
& \int_{0}^{\infty} e^{-s} s^{\mu} L_{m}^{(\mu)}(s) L_{k}^{(\mu)}(s) d s=\Gamma(\mu+1)\binom{m+\mu}{m} \delta_{m k} \\
& \int_{-\infty}^{\infty} e^{-s^{2}} H_{m}(s) H_{k}(s) d s=\pi^{\frac{1}{2}} 2^{m} m!\delta_{m k}
\end{aligned}
$$

And they are complete in $L^{2}\left(\mathbf{R}^{+}, e^{-s} s^{\mu} d s\right)$ and $L^{2}\left(\mathbf{R}, e^{-s^{2}} d s\right)$ respectively (cf [10], Chapter V ). We also need following notations. For

$$
\alpha=\sum_{j=1}^{r} \alpha_{j} c_{j}+\sum_{j<k} \sum_{l=1}^{d} \alpha_{j k}^{l} c_{j k}^{l} \in \mathcal{A},
$$

we set

$$
\begin{aligned}
& |\alpha|=\sum_{j=1}^{r} \alpha_{j}+\sum_{j<k} \sum_{l=1}^{d} \alpha_{j k}^{l}, \\
& \alpha!=\prod_{j=1}^{r} \alpha_{j}!\prod_{j<k} \prod_{l=1}^{d} \alpha_{j k}^{l}!, \\
& \alpha^{0}=\left(\alpha_{1}, \cdots, \alpha_{r}\right), \\
& \left|\alpha^{0}\right|=\sum_{j=1}^{r} \alpha_{j}, \\
& \alpha^{0}!=\prod_{j=1}^{r} \alpha_{j}!.
\end{aligned}
$$

For $\mathbf{s} \in \mathbf{C}^{r}, \lambda \in \mathbf{C}$, we write

$$
\mathbf{s}+\lambda=\left(s_{1}+\lambda, \cdots, s_{r}+\lambda\right)
$$

The gamma function of the symmetric cone $\Omega$ is defined by

$$
\begin{aligned}
\Gamma_{\Omega}(\mathbf{s}) & =\int_{\Omega} e^{-\operatorname{tr}(x)} \Delta_{\mathbf{s}}(x) \Delta(x)^{-\frac{n}{r}} d x \\
& =(2 \pi)^{\frac{n-r}{2}} \prod_{j=1}^{r} \Gamma\left(s_{j}-(j-1) \frac{d}{2}\right), \quad \operatorname{Re} s_{j}>(j-1) \frac{d}{2}, j=1, \cdots, r .
\end{aligned}
$$

Assume that

$$
v=\sum_{j=1}^{r} v_{j} c_{j}+\sum_{j<k} \sum_{l=1}^{d} v_{j k}^{l} c_{j k}^{l} \in V_{+},
$$

we define

$$
\psi_{\alpha}(v)=C \prod_{j=1}^{r} e^{-v_{j}^{2}} v_{j}^{\nu-\frac{n}{r}} L_{\alpha_{j}}^{\left(\mu_{j}\right)}\left(2 v_{j}^{2}\right) \prod_{j<k} \prod_{l=1}^{d} e^{-\frac{1}{2}\left(v_{j k}^{l}\right)^{2}} H_{\alpha_{j k}^{l}}\left(v_{j k}^{l}\right)
$$

where

$$
\begin{aligned}
& C=2^{\frac{1}{2}\left(\nu r-n+|\alpha|-\left|\alpha^{0}\right|\right)} \alpha^{0}!(\alpha!)^{-\frac{1}{2}} \Gamma_{\Omega}\left(\alpha^{0}+\nu-\frac{n}{r}\right)^{-\frac{1}{2}}, \\
& \mu_{j}=\nu-\frac{n}{r}-1-\frac{d}{2}(j-1), \quad \nu>1+d(r-1) .
\end{aligned}
$$

We regard $\psi_{\alpha}(v)$ as the functions on group $T$. From the orthogonal relations and completeness of the Laguerre polynomials and Hermite polynomials, we conclude that $\left\{\psi_{\alpha}(v): \alpha \in \mathcal{A}\right\}$ is an orthonormal basis of $L^{2}\left(T, d \mu_{l}\right)$.

Now we define a family of functions $\phi_{\alpha}^{\varepsilon}$ by

$$
\hat{\phi}_{\alpha}^{\varepsilon}(y)= \begin{cases}\psi_{\alpha}(v), & y=t(v)^{\prime} P(\varepsilon) e, \\ 0, & y \notin \Omega_{\varepsilon} .\end{cases}
$$

Then $\phi_{\alpha}^{\varepsilon} \in H_{\varepsilon}$ is admissible and $\left\{\phi_{\alpha}^{\varepsilon}: \varepsilon \in E, \alpha \in \mathcal{A}\right\}$ is an complete orthonormal system with respect to $\langle\cdot, \cdot\rangle_{A W}$.

## 4. The decomposition of $L_{\nu}^{2}\left(T_{\Omega}\right)$

The weighted $L^{2}$-space on the symmetric tube domain $T_{\Omega}$ is defined by

$$
L_{\nu}^{2}\left(T_{\Omega}\right)=\left\{F:\|F\|_{\nu}^{2}=\int_{T_{\Omega}}|F(x+i y)|^{2} \Delta(y)^{\nu-\frac{2 n}{r}} d x d y<\infty\right\} .
$$

$\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)$, the weighted Bergman space, is the subspace of all holomorphic functions in $L_{\nu}^{2}\left(T_{\Omega}\right)$, i.e.,

$$
\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)=\left\{F \in L_{\nu}^{2}\left(T_{\Omega}\right): F \text { is holomorphic on } T_{\Omega}\right\} .
$$

We assume that $\nu>1+d(r-1)$ so that $\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right) \neq\{0\}$. We want to decompose $L_{\nu}^{2}\left(T_{\Omega}\right)$ into the direct sum of subspaces such that the first component is exactly the weighted Bergman space $\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)$.

Suppose $f \in H_{\varepsilon}, \phi \in A W_{\varepsilon}$. For $\nu>1+d(r-1)$, we define the weighted wavelet transform $W_{\phi}^{\nu}$ by

$$
F(u+i t(v) e)=W_{\phi}^{\nu} f(u+i t(v) e)=C_{\phi}^{-\frac{1}{2}} W_{\phi} f(u, v) \Delta(t(v) e)^{-\frac{\nu}{2}} .
$$

Set

$$
\mathcal{H}_{\alpha}^{\varepsilon}=\left\{F=W_{\phi_{\alpha}^{\varepsilon}}^{\nu} f: f \in H_{\varepsilon}\right\} .
$$

Proposition 2. For $\nu>1+d(r-1)$, we have

$$
\begin{equation*}
L_{\nu}^{2}\left(T_{\Omega}\right)=\bigoplus_{\varepsilon \in E, \alpha \in \mathcal{A}} \mathcal{H}_{\alpha}^{\varepsilon} \tag{4.1}
\end{equation*}
$$

and

$$
\mathcal{H}_{0}^{e}=\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)
$$

Proof. By Theorem 2, $W_{\phi}^{\nu}$ is an isometric operator from $H_{\varepsilon}$ into $L_{\nu}^{2}\left(T_{\Omega}\right)$ and $\mathcal{H}_{\alpha}^{\varepsilon}$ 's are mutually orthogonal subspaces of $L_{\nu}^{2}\left(T_{\Omega}\right)$. We need to prove

$$
L_{\nu}^{2}\left(T_{\Omega}\right) \subset \bigoplus_{\varepsilon \in E, \alpha \in \mathcal{A}} \mathcal{H}_{\alpha}^{\varepsilon}
$$

Suppose $F \in L_{\nu}^{2}\left(T_{\Omega}\right)$. Write $F_{v}(u)=F(u+i t(v) e)$. For $v \in V_{+}$almost every where, $F_{v} \in L^{2}(V)$. We let $G(v, y)=\hat{F}_{v}(y)$. Fix $y=t\left(v^{1}\right)^{\prime} P(\varepsilon) e \in \Omega_{\varepsilon}$, then $G(v, y)$, regarded as the function on $T$, is in $L^{2}\left(T, \Delta(t(v) e)^{\nu-\frac{n}{r}} d \mu_{l}(t(v))\right)$. Since

$$
\left\{\Delta(t(v) e)^{-\frac{\nu}{2}+\frac{n}{2 r}} \hat{\phi}_{\alpha}^{\varepsilon}\left(t(v)^{\prime} y\right): \alpha \in \mathcal{A}\right\}
$$

is an orthonormal basis of $L^{2}\left(T, \Delta(t(v) e)^{\nu-\frac{n}{r}} d \mu_{l}(t(v))\right)$, we get

$$
G(v, y)=\sum_{\alpha \in \mathcal{A}} a_{\alpha}(y) \Delta(t(v) e)^{-\frac{\nu}{2}+\frac{n}{2 r}} \hat{\phi}_{\alpha}^{\varepsilon}\left(t(v)^{\prime} y\right), \quad y \in \Omega_{\varepsilon} .
$$

We define the functions $f_{\alpha}^{\varepsilon}$ by

$$
\hat{f}_{\alpha}^{\varepsilon}(y)= \begin{cases}a_{\alpha}(y), & y \in \Omega_{\varepsilon} \\ 0, & y \notin \Omega_{\varepsilon}\end{cases}
$$

It is easy to see that $f_{\alpha}^{\varepsilon} \in H_{\varepsilon}$. Therefore we have

$$
F(u+i t(v) e)=\sum_{\varepsilon \in E, \alpha \in \mathcal{A}} W_{\phi_{\alpha}^{\varepsilon}}^{\nu} f_{\alpha}^{\varepsilon}(u+i t(v) e) .
$$

This proves (4.1).
Now we assume that

$$
F(u+i t(v) e)=W_{\phi_{0}^{e}}^{\nu} f(u+i t(v) e), \quad f \in H_{e}
$$

Note that

$$
\begin{aligned}
& \operatorname{tr}(y)=\sum_{j=1}^{r} y_{j}=\sum_{j=1}^{r} v_{j}^{2}+\frac{1}{2} \sum_{j<k} \sum_{l=1}^{d}\left|v_{j k}^{l}\right|^{2}, \\
& \Delta(y)=\prod_{j=1}^{r} v_{j}^{2}, \quad y=t(v)^{\prime} e,
\end{aligned}
$$

we have

$$
\hat{\phi}_{0}^{e}(y)= \begin{cases}\Gamma_{\Omega}\left(\nu-\frac{n}{r}\right)^{-\frac{1}{2}} 2^{\frac{\nu r-n}{2}} e^{-\operatorname{tr} y} \Delta(y)^{\frac{\nu}{2}-\frac{n}{2 r}}, & y \in \Omega \\ 0, & y \notin \Omega\end{cases}
$$

Therefore,

$$
\begin{aligned}
F(u+i t(v) e) & =\Delta(t(v) e)^{-\frac{\nu}{2}-\frac{n}{2 r}} \int_{V} f(x) \overline{\phi\left(t(v)^{-1} x-t(v)^{-1} u\right)} d x \\
& =(2 \pi)^{-n} \Delta(t(v) e)^{-\frac{\nu}{2}+\frac{n}{2 r}} \int_{\Omega} \hat{f}(y) e^{i(u \mid y)} \overline{\hat{\phi}\left(t(v)^{\prime} y\right)} d y \\
& =(2 \pi)^{-n} \Gamma_{\Omega}\left(\nu-\frac{n}{r}\right)^{-\frac{1}{2}} \int_{\Omega} e^{i(u+i t(v) e \mid y)} \hat{f}(y) \Delta(2 y)^{\frac{\nu}{2}-\frac{n}{2 r}} d y
\end{aligned}
$$

where we make use of the equality

$$
\Delta(t(v) e)=\operatorname{Det}(t(v))^{\frac{r}{n}}=\operatorname{Det}\left(t(v)^{\prime}\right)^{\frac{r}{n}}=\Delta\left(t(v)^{\prime} e\right)
$$

We define the map $\mathcal{F}_{\nu}$ by

$$
g(y)=\mathcal{F}_{\nu} f(y)=(2 \pi)^{-\frac{n}{2}} \Gamma_{\Omega}\left(\nu-\frac{n}{r}\right)^{-\frac{1}{2}} \hat{f}(y) \Delta(2 y)^{-\frac{\nu}{2}+\frac{n}{2 r}}
$$

and the map $\mathcal{L}_{\nu}$ by

$$
F(u+i t(v) e)=\mathcal{L}_{\nu} g(u+i t(v) e)=(2 \pi)^{-\frac{n}{2}} \int_{\Omega} e^{i(u+i t(v) e \mid y)} g(y) \Delta(2 y)^{\nu-\frac{n}{r}} d y
$$

It is obvious that $\mathcal{F}_{\nu}$ is an isomorphism from $H^{2}(V)$ onto $L_{\nu}^{2}(\Omega)=$ $L^{2}\left(\Omega, \Delta(2 y)^{\nu-\frac{n}{r}} d y\right) . \mathcal{L}_{\nu}$ is an isomorphism from $L_{\nu}^{2}(\Omega)$ onto $\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)($ cf $[3])$. We see that $W_{\phi_{0}^{e}}^{\nu}=\mathcal{L}_{\nu} \circ \mathcal{F}_{\nu}$ and

$$
\|F\|_{L_{\nu}^{2}\left(T_{\Omega}\right)}^{2}=\Gamma_{\Omega}\left(\nu-\frac{n}{r}\right)\|g\|_{L_{\nu}^{2}(\Omega)}^{2}=\|f\|_{L^{2}(V)}^{2} .
$$

Clearly, $\mathcal{H}_{0}^{e}$ is exactly the weighted Bergman space $\mathcal{H}_{\nu}^{2}\left(T_{\Omega}\right)$.
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## References

[1] Bernier, D., and K. F. Taylor, Wavelets from square-integrable representations, SIAM J. Math. Anal. 27 (1996), 594-608.
[2] Duflo, M., and C. C. Moore, On the regular representation of a nonunimodular locally compact group, J. Funct. Anal. 21 (1976), 209-243.
[3] Faraut, J., and A. Korányi, "Analysis on symmetric cones," Oxford, 1994.
[4] Grossmann, A., and J. Morlet, Decomposition of Hardy functions into square integrable wavelets of constant shape, SIAM J. Math. Anal. 15 (1984), 723-736.
[5] Grossmann, A., J. Morlet, and T. Paul, Transforms associated to square integrable group representations I: General results, J. Math. Phys. 26 (1985), 2473-2479.
[6] Heil, C. E., and D. F. Walnut, Continuous and discrete wavelet transforms, SIAM Review 31 (1989), 628-666.
[7] Jiang, Q., and L. Peng, Wavelet transform and Toeplitz-Hankel type operators, Math. Scand. 70 (1992), 247-264.
[8] Liu, H., and L. Peng, Admissible wavelets associated with the Heisenberg group, Pacific J. Math. 180 (1997), 101-123.
[9] Paul, T., Functions analytic on the half-plane as quantum mechanical states, J. Math. Phys. 25 (1984), 3252-3263.
[10] Szegö, G., "Orthogonal polynomials," Amer. Math. Soc. Colloq. Publications, Vol. 23, Revised edition, 1959.
[11] Terras, A., "Harmonic analysis on symmetric spaces and applications II," Springer-Verlag, 1988.

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